

Conservation Properties of the Bridging Domain Method for Coupled Molecular/Continuum Dynamics

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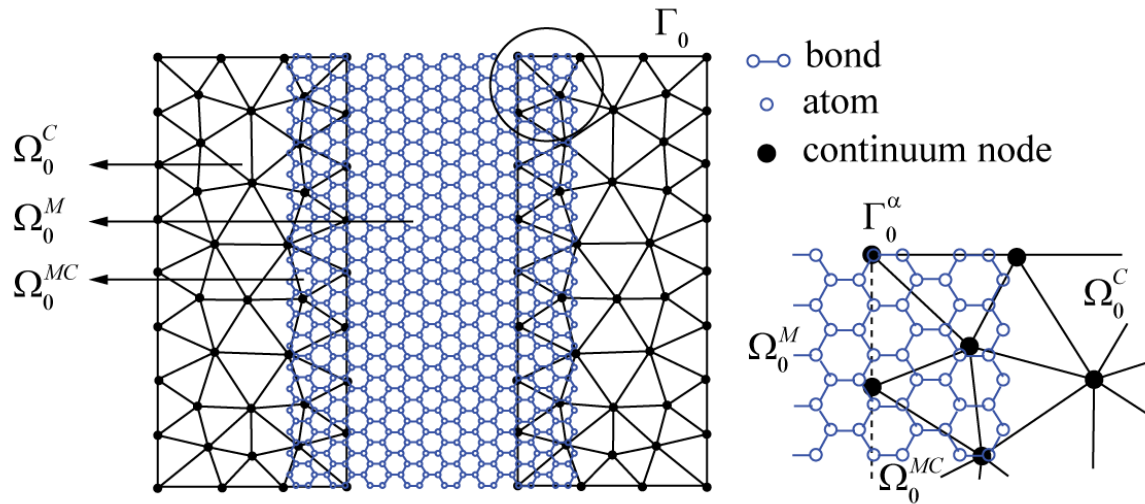
Theoretical and Applied Mechanics



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- Bridging domain method
- Conservation properties
 - Linear momentum
 - Angular momentum
 - Energy
- Numerical tests
 - One-dimensional
 - Two-dimensional
- Conclusions



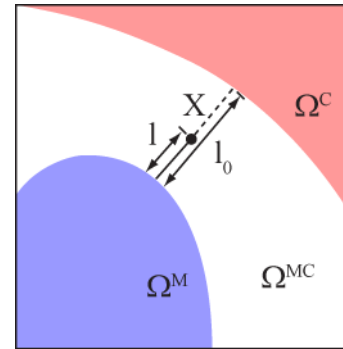
■ Main features

- Couples atomistic and continuum model
- Scaling factor function is introduced to the coupling region
- Total energy of the system is the combination of the energy of the atomistic and continuum model
- Lagrange multipliers are used to enforce the compatibility between the models



- Scaling factor (linear)

$$\alpha(\mathbf{X}) = \begin{cases} 0 & \text{if } \mathbf{X} \in \Omega_0^M / \Omega_0^{MC} \\ l(\mathbf{X}) / l_0(\mathbf{X}) & \text{if } \mathbf{X} \in \Omega_0^{MC} \\ 1 & \text{if } \mathbf{X} \in \Omega_0^C / \Omega_0^{MC} \end{cases}$$



- Lagrange multiplier

- Interpolated from a continuum field

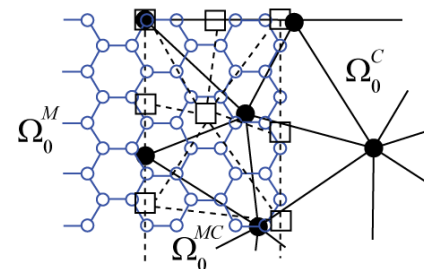
$$W = W^{\text{int}} - W^{\text{ext}} + \lambda^T \mathbf{g}$$

$$g_{ij} = \sum_J N_J(\mathbf{X}_I) u_{ij} - d_{ij}$$

N_{JI}

$$\lambda(\mathbf{X}_I) = \sum_J N_{JI}^\lambda \bar{\lambda}_J$$

□ Lagrange multiplier node





Equations of motion

$$\bar{M}_I^M \ddot{d}_{il} + \bar{f}_{il}^{\text{int}M} + \bar{f}_{il}^{GM} = 0$$

$$\bar{M}_I^M \ddot{u}_{il} + \bar{f}_{il}^{\text{int}C} + \bar{f}_{il}^{GC} = 0 \quad \text{where}$$

$$\sum_{I \in S} N_{IJ} u_{il} - d_{ij} = 0$$

$$\bar{f}_{il}^{\text{int}M} = \sum_{J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{d_{il}}$$

$$\bar{f}_{il}^{GM} = \frac{\partial G}{\partial d_{il}} = -\lambda_{il}$$

$$\bar{f}_{il}^{\text{int}C} = \int_{\Omega_0^C} \alpha \frac{\partial N_I}{\partial X_j} P_{ji} d\Omega_0^C$$

$$\bar{f}_{il}^{GC} = \frac{\partial G}{\partial u_{il}} = \sum_{J \in M} \lambda_{iJ} N_{IJ}$$

Conservation properties

- Linear momentum

$$\mathbf{p} = \sum_{I \in M} \bar{M}_I^M \dot{\mathbf{d}}_I + \sum_{I \in S} \bar{M}_I^C \dot{\mathbf{u}}_I$$

- Angular momentum

$$\mathbf{L} = \sum_{I \in M} \bar{M}_I^M \mathbf{x}_I^M \times \dot{\mathbf{d}}_I + \sum_{I \in S} \bar{M}_I^C \mathbf{x}_I^C \times \dot{\mathbf{u}}_I$$

- Energy

$$E = \frac{1}{2} \sum_{I \in M} \bar{M}_I^M \dot{d}_{il} \dot{d}_{il} + \frac{1}{2} \sum_{I \in S} \bar{M}_I^C \dot{u}_{il} \dot{u}_{il} + \sum_{I < J \in M} \bar{w}_M(\mathbf{x}_I, \mathbf{x}_J) + \int_{\Omega_0^C} \bar{w}_C(\mathbf{F}) d\Omega_0^C$$



Conservation of Linear Momentum

$$\begin{aligned} \frac{dp_i}{dt} &= \frac{d}{dt} \left(\sum_{I \in M} \bar{M}_I^M \dot{d}_{iI} + \sum_{I \in S} \bar{M}_I^C \dot{u}_{iI} \right) \\ &= \sum_{I \in M} \lambda_{iI} - \sum_{I, J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{\partial d_{iI}} - \sum_{I \in M} \lambda_{iI} \sum_{J \in S} N_{JI} - \sum_{I \in S} \int_{\Omega_0^C} \alpha \frac{\partial N_I}{\partial X_i} P_{ji} d\Omega_0^C \end{aligned}$$

↑ cancel ↑
↑ vanish ↑

$$\sum_{J \in M} N_J(\mathbf{X}) = 1 \quad \text{partition of unity}$$

$$\sum_{I, J \in M} \frac{\partial w_{IJ}^M}{\partial d_{iI}} = \frac{1}{2} \sum_{I, J \in M} \left(\frac{\partial w_{IJ}^M}{\partial d_{iI}} + \frac{\partial w_{IJ}^M}{\partial d_{iJ}} \right) = 0 \quad \text{anti-symmetry}$$



Conservation of Angular Momentum

$$\frac{d\mathbf{L}}{dt} = \frac{d}{dt} \left(\sum_{I \in M} \bar{M}_I^M \mathbf{x}_I^M \times \dot{\mathbf{d}}_I + \sum_{I \in S} \bar{M}_I^C \mathbf{x}_I^C \times \dot{\mathbf{u}}_I \right) \Rightarrow \frac{dL_k}{dt} = \varepsilon_{jik} \left(\sum_{I \in M} \bar{M}_I^M x_{jl}^M \ddot{d}_{il} + \sum_{I \in S} \bar{M}_I^C x_{jl}^C \ddot{u}_{il} \right)$$

$$\frac{dL_k}{dt} = \varepsilon_{jik} \left(\sum_{I \in M} \left[\lambda_{iI} - \sum_{J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{\partial d_{il}} \right] x_{jl}^M + \sum_{I \in S} \left[\sum_{J \in M} \lambda_{iJ} N_{IJ} - \int_{\Omega_0^C} \alpha \frac{\partial N_I}{\partial X_l} P_{li} d\Omega_0^C \right] x_{jl}^C \right)$$

cancel

F_{jl}

$$\sum_{I \in S} \varepsilon_{jik} F_{jl} P_{li} = \frac{1}{2} \sum_{I \in S} \varepsilon_{jik} (F_{jl} P_{li} - F_{il} P_{lj}) = \frac{1}{2} \sum_{I \in S} \varepsilon_{jik} J (\sigma_{ji} - \sigma_{ij}) = 0 \quad \text{Cauchy stress is symmetric}$$

$$\varepsilon_{jik} \sum_{J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{\partial d_{il}} x_{jl}^M = \frac{1}{2} \varepsilon_{jik} \sum_{J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{\partial d_{il}} (x_{jl}^M - x_{jJ}^M) = \frac{1}{2} \sum_{J \in M} (1 - \alpha_{IJ}) \mathbf{r}_{IJ} \times \frac{\partial w_{IJ}^M}{\partial \mathbf{d}_I} = 0$$

Bond force is co-axial with the bond vector



Conservation of Energy

$$\frac{dE}{dt} = \sum_{I \in M} \bar{M}_I^M \ddot{d}_{il} \dot{d}_{il} + \sum_{I \in S} \bar{M}_I^C \ddot{u}_{il} \dot{u}_{il} + \frac{1}{2} \sum_{I, J \in M} (1 - \alpha_{IJ}) \left(\frac{\partial w_{IJ}^M}{\partial d_{il}} \dot{d}_{il} + \frac{\partial w_{IJ}^M}{\partial d_{ij}} \dot{d}_{ij} \right) + \int_{\Omega_0^C} \alpha \frac{\partial w^C}{\partial F_{ij}} \dot{F}_{ij} d\Omega_0^C$$

$$\frac{dE}{dt} = - \sum_{I \in M} (\bar{f}_{il}^M + \bar{f}_{il}^{GM}) \dot{d}_{il} - \sum_{I \in S} (\bar{f}_{il}^{GC} + \bar{f}_{il}^C) \dot{u}_{il} + \sum_{J \in M} (1 - \alpha_{IJ}) \frac{\partial w_{IJ}^M}{\partial d_{il}} \dot{d}_{il} + \int_{\Omega_0^C} \alpha P_{ji} \frac{\partial N_I}{\partial X_j} \dot{u}_{il} d\Omega_0^C$$

cancel
cancel

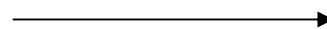
$$\frac{dE}{dt} = - \sum_{I \in M} \bar{f}_{il}^{GM} \dot{d}_{il} - \sum_{I \in S} \bar{f}_{il}^{GC} \dot{u}_{il} = \dot{G} \quad \Rightarrow \text{always zero}$$



■ Constraint Matrix

$$\sum_{J \in M} \lambda_{IJ} A_{IJ} = B_{ii}$$

$$A_{IJ} = -\sum_{L \in S} \frac{N_{LI} N_{LJ}}{\bar{M}_L^C} - \frac{\delta_{IJ}}{\bar{M}_I^M}$$

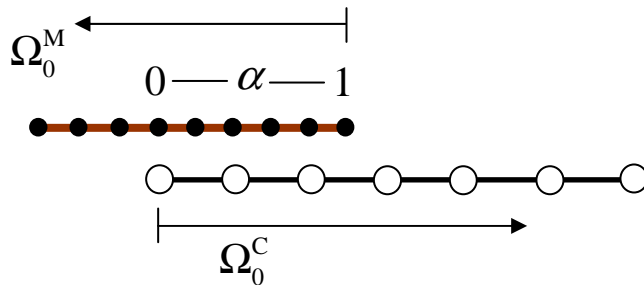


$$A_{II} = \sum_{J \in S} A_{IJ}$$

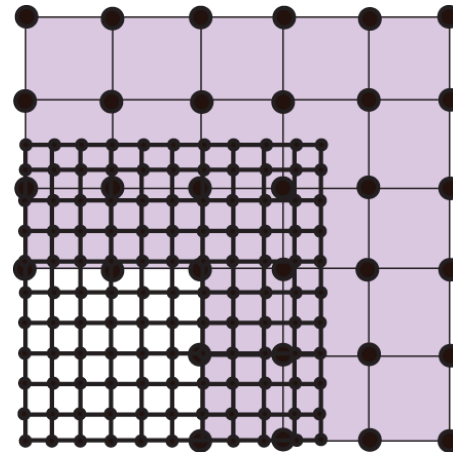
consistent constraint matrix
-compatibility is satisfied exactly
-computationally more expensive

diagonalized constraint matrix
-compatibility is weakly enforced
-cheaper in computation

■ Numerical Models



1D bridging domain model
(Lennard-Jones potential)



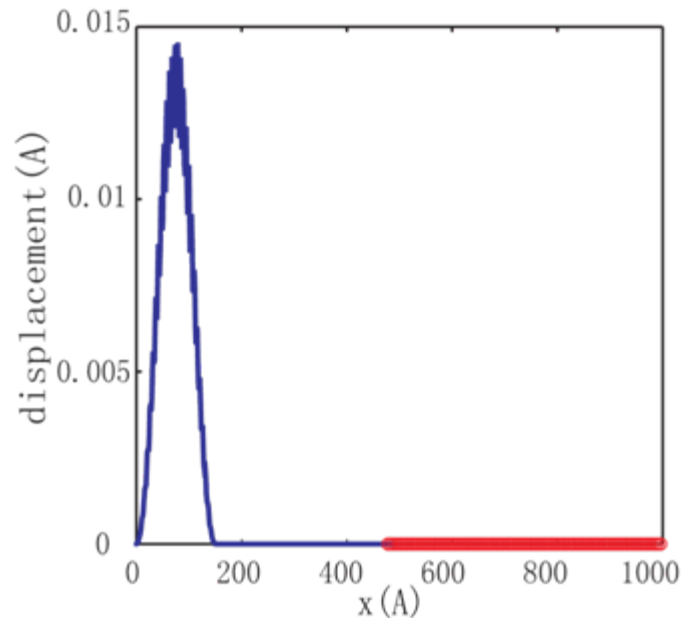
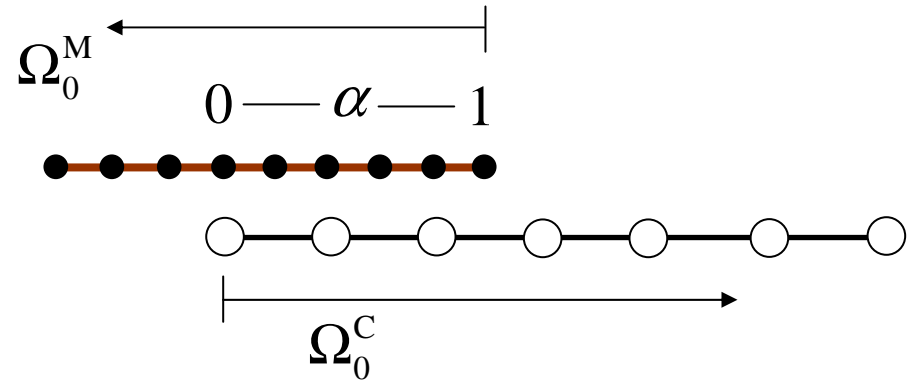
2D bridging domain model
(Quadratic potentials)



- One-dimensional example
 - Lennard-Jones potential

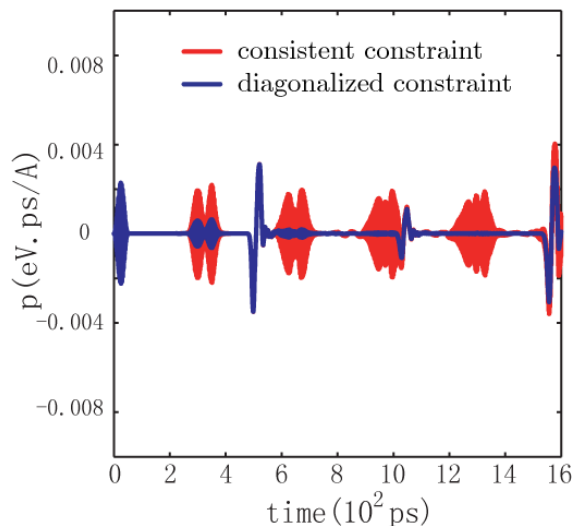
$$w(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right]$$

- Initial condition

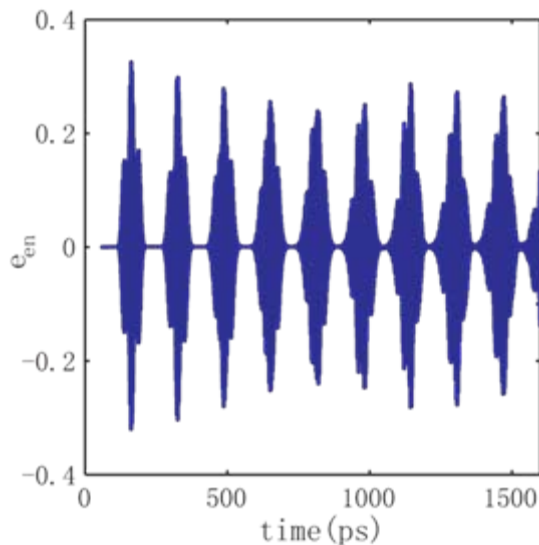




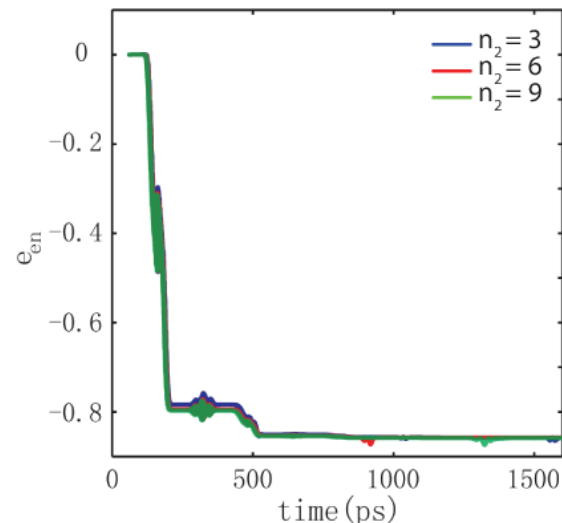
Numerical Test – 1D



linear momentum history



error in energy with
consistent constraint

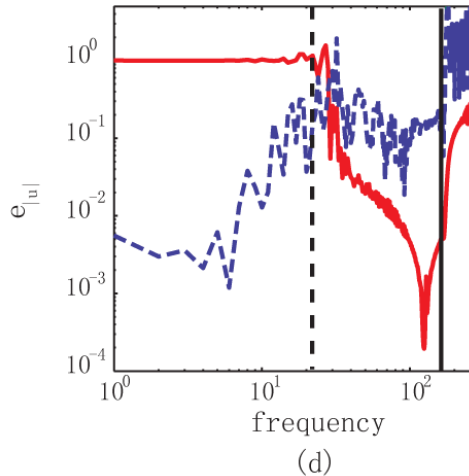
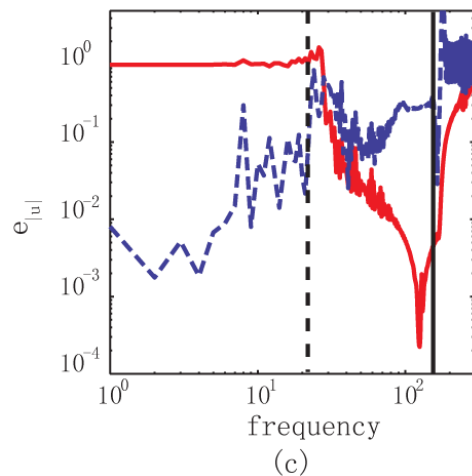
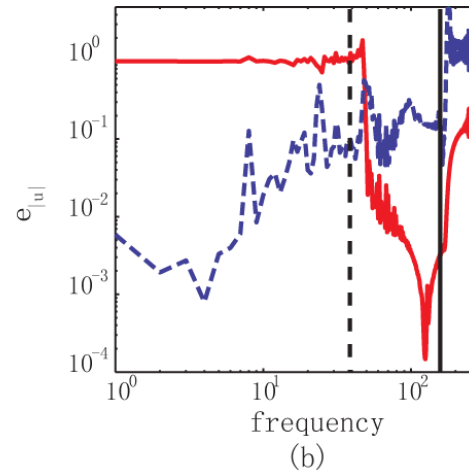
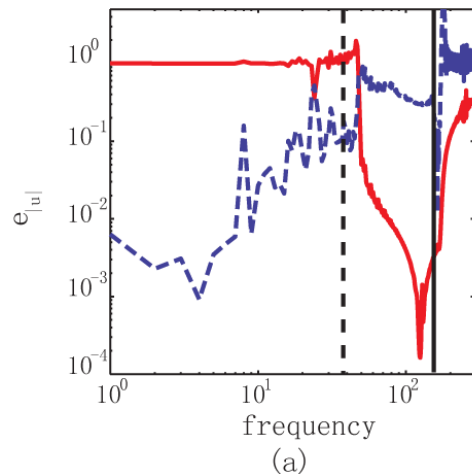


error in energy with
diagonalized constraint

$$error = \frac{A(t) - A_0}{A_0}$$



Spurious Reflection

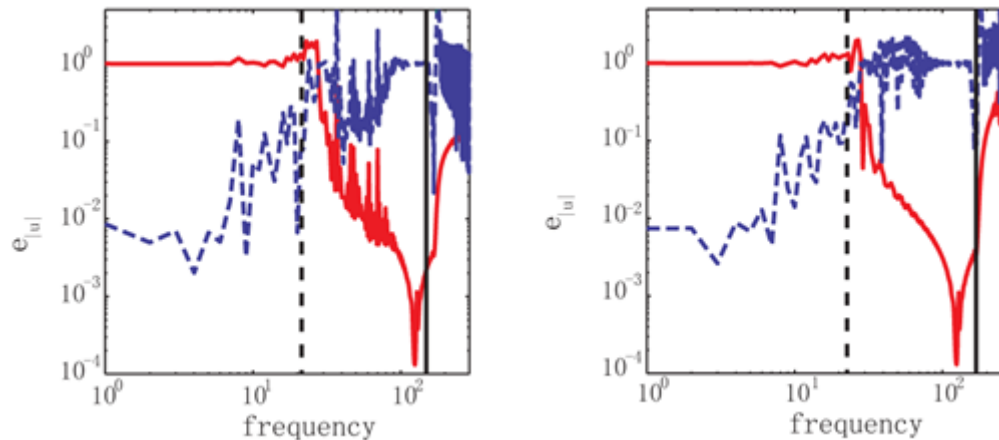


— transmission
- - - reflection

Reflectivity and transmission for n_2 overlaid elements with n_1 atoms per element with *diagonalized constraint matrix*. (a) $n_1 = 4, n_2 = 3$
(b) $n_1 = 4, n_2 = 6$ (c) $n_1 = 7, n_2 = 3$
(d) $n_1 = 7, n_2 = 6$



■ Spurious Reflection



Reflectivity and transmission for the **consistent constraint matrix**. (a) $n_1 = 7, n_2 = 3$ (b) $n_1 = 7, n_2 = 6$

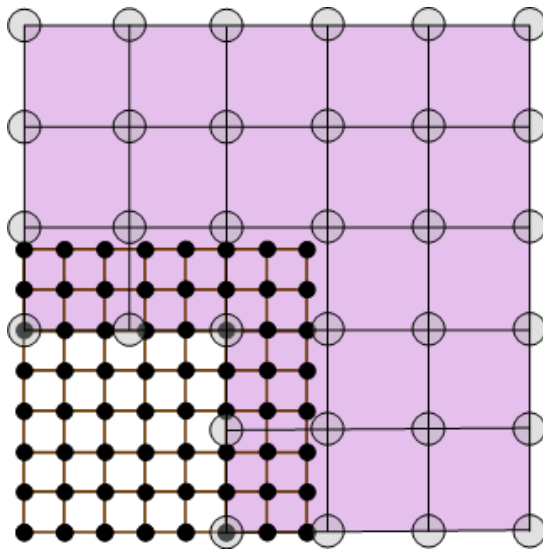
- *Low frequencies smoothly propagate into coarser scale.*
- *The diagonalized constraint matrix can effectively eliminate the spurious reflection at the atomistic/continuum interface.*
- *One-dimensional examples show that bigger overlapping domain gives better transmission of the low frequency wave.*



■ Two-dimensional example

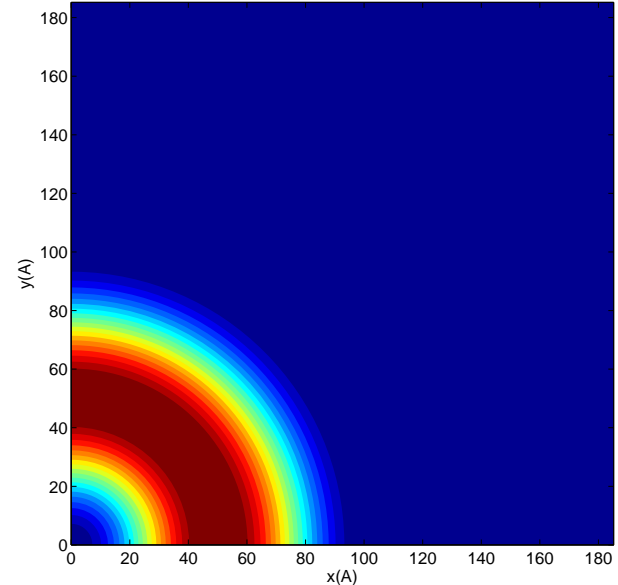
- Quadratic potential

$$U(r, \theta) = \frac{1}{2} k_a (r - r_0)^2 + \frac{1}{2} k_\theta (\theta - \theta_0)^2$$



- atomic node
- continuum node

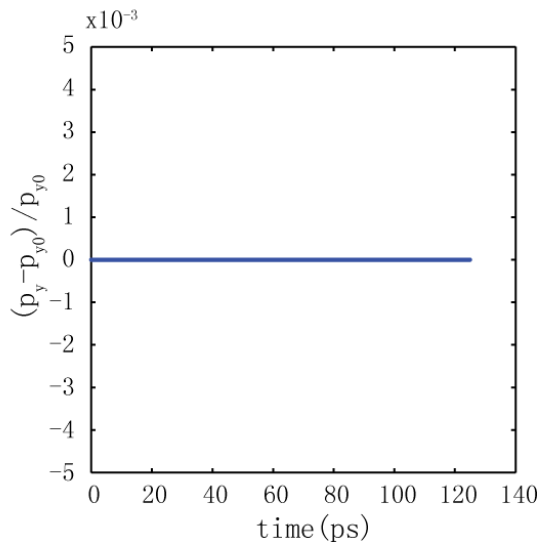
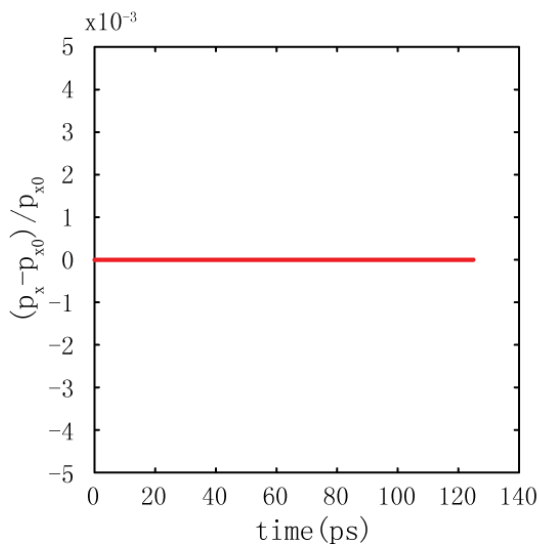
Coupling model



Initial displacement field

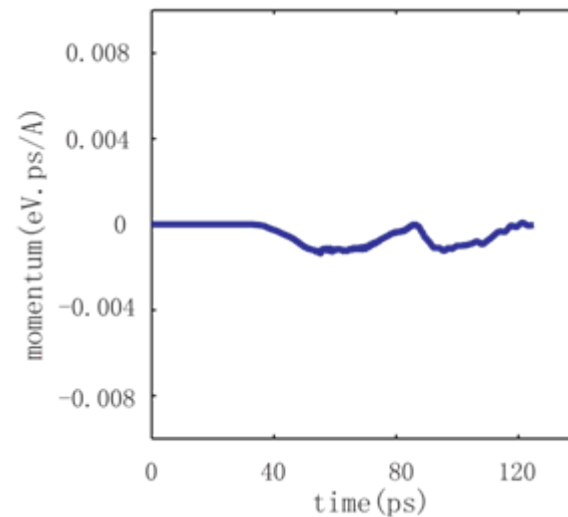


Numerical Test – 2D



error in linear momentum of 2D tests

$$n_1 = 4, n_2 = 1$$

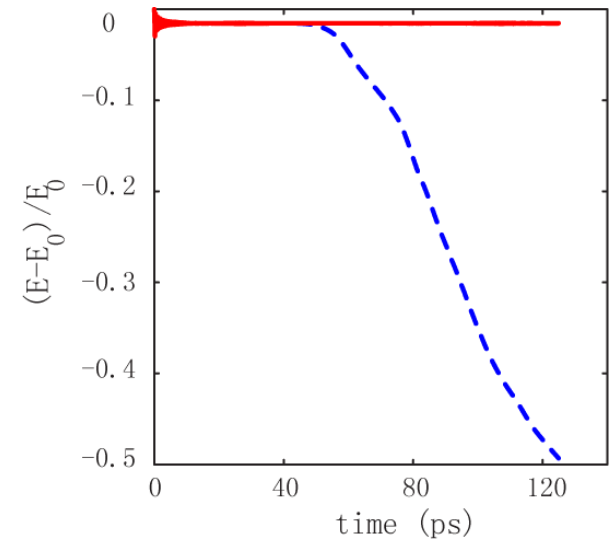


time history angular
momentum of 2D tests



Numerical Test – 2D

- *Both linear momentum and angular momentum are conserved.*
- *Energy conservation depends on the compatibility condition between the atomistic and continuum models. The consistent constraint matrix conserves energy while the diagonalized constraint matrix dissipates energy.*



error in energy of
two-dimensional tests



Conclusions

- The bridging domain method conserves energy and momentum
- Oscillation of linear momentum and energy (1D), and angular momentum (2D) is observed as waves pass through the coupling domain
 - These quantities are restored after the wave exits the coupling domain
- Energy conservation depends on the constraint implementation:
 - The consistent constraint conserves total energy
 - The diagonalized constraint dissipates high frequency energy, eliminating reflections at the interface between different scales